Hartree–Fock Theory Implementation in Python

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1 Introduction

Hartree–Fock (HF) method is a method of approximation for the determination of the wave function and the energy of a chemical system. The Hartree–Fock method assumes that the exact N-body wave function of the system can be approximated by a single Slater determinant or by a single permanent N spin-orbitals. By invoking the variational method, one can derive a set of N-coupled equations for the N spin orbitals. A solution of these equations yields the Hartree–Fock wave function and energy of the system. In this report, I will construct the Self-Consistent-Field (SCF) procedure to carry out Hartree–Fock calculations.

2 SCF Procedure

The main objective of this report is to implement the SCF procedure. In order to achieve that, we will go through these steps:

- 1. Specify molecule, basis functions.
- 2. Form the overlab matrix **S** from the basis functions.
- 3. Provide an initial guess for the orbitals C.
- 4. Form the density matrix $\mathbf{D_{pq}} = \sum_{i}^{occ} \mathbf{c_{pi}^* c_{qi}}$.
- 5. Form the Fock matrix $\mathbf{F} = \mathbf{H} + 2\mathbf{J} \mathbf{K}$
- 6. Solve the Hartree-Fock-Roothan equation: FC = SCE to get a new orbital matrix C.
- 7. Calculate the Fock matrix **F** using the new orbital matrix **C**.
- 8. Repeat from step 4 until the orbital matrix C and the orbital energies converges.

Now that we know how to design the SCF procedure, we are left to compute all the necessary mathematical objects which will be discussed in the next section.

3 Computations

In order to design the SCF, we will need to compute all of the matrices required to satisfy the Hartree–Fock–Roothan (HFR) equation:

$$\sum_{\nu} F_{\mu\nu} C_{\nu i} = \sum_{\nu} S_{\mu\nu} C_{\nu i} E_i$$
$$FC = SCE$$

Therefore, we shall compute all of the following:

- 1. The Fock matrix \mathbf{F} .
- 2. The matrix of wavefunction amplitudes for each MO C.
- 3. The overlab matrix S.
- 4. The density matrix **D**.

3.1 The Overlap Matrix S

3.1.1 Constructing S

The overlap matrix \mathbf{S} contains the information about how the basis functions ϕ overlaps. Each element in the matrix is the Bra-ket of those basis functions. We shall form it using a matrix \mathbf{B} that contains the basis vectors as its columns:

$$S_{ij} = \langle \phi_i | \phi_j \rangle$$

$$B = \begin{pmatrix} | & | & | \\ |\phi_1 \rangle & |\phi_2 \rangle & |\phi_3 \rangle \\ | & | & | \end{pmatrix} \quad B^{\dagger} = \begin{pmatrix} - & \langle \phi_1 | & - \\ - & \langle \phi_2 | & - \\ - & \langle \phi_3 | & - \end{pmatrix}$$

$$B^{\dagger}B = \begin{pmatrix} - & \langle \phi_1 | & - \\ - & \langle \phi_2 | & - \\ - & \langle \phi_2 | & - \\ - & \langle \phi_3 | & - \end{pmatrix} \begin{pmatrix} | & | & | \\ |\phi_1 \rangle & |\phi_2 \rangle & |\phi_3 \rangle \\ | & | & | \end{pmatrix} \equiv S = \begin{pmatrix} \langle \phi_1 | \phi_1 \rangle & \langle \phi_1 | \phi_2 \rangle & \langle \phi_1 | \phi_3 \rangle \\ \langle \phi_2 | \phi_1 \rangle & \langle \phi_2 | \phi_2 \rangle & \langle \phi_2 | \phi_3 \rangle \\ | & | & | & | \end{pmatrix}$$

If we have an orthonormal set of basis functions ϕ_i , then the overlap matrix S will be the same as the identity matrix. Since the diagonal elements represents the normalization $\langle \phi_i | \phi_i \rangle = 1$ and the off-diagonal element will represent the orthogonality $\langle \phi_i | \phi_j \rangle = 0, i \neq j$. However, our basis functions are not necessarly orthonormal, so we will need to **orthogonalize** S into an orthonormal representation.

3.1.2 Orthogonalizing S

In order to orthogonalize S, we will use some Linear Algebra to construct an orthogonalizing matrix A. S is summetric, and we can choose it to be real. Since S is symmetric and real, A is also symmetric and real. Therfore $A^{\dagger} = A$. This matrix is important and we will use it in thext articles.

$$A^{\dagger}SA = 1$$

$$A^{\dagger}SA = ASA = A^{2}S = 1$$

$$\Rightarrow A = S^{-1/2}$$

Now that we found A, we will use it to transform $S \to S'$, S' is S represented in the orthonormal basis set. We will do that by transforming B and B^{\dagger} using A:

$$B' = BA$$

$$S' = B'^{\dagger}B'$$

$$S' = (BA)^{\dagger}(BA) = A^{\dagger}B^{\dagger}BA$$

$$S' = A^{\dagger}SA = ASA$$

3.2 The Fock Matrix F

3.2.1 Transforming $F, C \rightarrow F', C'$

Now we have S represented in the orthonormal basis set. We will carry out the same transofmation for the rest of HFR equation:

$$FC = SCE$$

$$F(\mathbf{1})C = S(\mathbf{1})CE$$

$$FAA^{-1}C = SAA^{-1}CE$$
 Mutliply by A from the left
$$AFAA^{-1}C = ASAA^{-1}CE$$

We can recognize S' = ASA, we can define F' = AFA and $C' = A^{-1}C$

$$F'C' = S'C'E$$
.

Since S' = 1 in the orthonormal basis set.

$$F'C' = C'E$$
.

3.2.2 Forming F

For closed-shell systems, F has the form F = H + 2J - K. Where H is the one-electron Hamiltonian H = T + V, J is the Coulomb integral matrix, and K is the exchange integral matrix. We will take H as our initial guess for F, then we need to solve HFR equation to get C. After obtaining C, we will use it to construct new F.

Now we will use C to obtain a new matrix D which is called **Density Matrix**, then we will compute a new tensor I which is called **Repulsive Tensor**. Using D and I we will have everything to form F.

$$D_{pq} = \sum_{i}^{occ} c_{pi}^* c_{qi}$$

Where $c_{pi}^*c_{qi}$ are the probability of some basis function p contributing to the MO i

$$I_{pqrs} = \int d\tau \; \phi_p^*(1)\phi_q(1) \frac{1}{r_{ij}} \phi_r^*(2)\phi_s(2)$$

Using D and I we can now get J and K

$$J_{pq} = \sum_{rs} D_{rs} I_{pqrs} \quad K_{ps} = \sum_{rq} D_{rq} I_{pqrs}$$

3.3 Finalizing SCF Procedure

Now we have everything we need, we will obtain the energy using this expression: $E = E_{nuc} + \sum_{pq} (H_{pq} + F_{pq}) D_{pq}$, where E_{nuc} is the nuclear repulsion energy.

By reaching here, we have completed one iteration of the SCF procedure. We can see that $F_{Guessed} \to C \to D, J, K \to F_{New} \to C \to D..., F$ (which depends on C) will yield a new C. This forms a Self-Consistent-Field SCF which we need to solve iteratively until our orbital coefficients and our orbital energies converge.

4 Results

After implementing the SCF procedure in Python, and using Psi4 library to compute all integrals. I tested it for four different molecules: HF, H_2O , NH_3 , CH_4 , and H_2S . I tested the code for two different basis sets: STO-3G and 6-31g(d). The accuracy was exact to the 10^8 th digit. However, there was some noticable time different in execution. Figure 1 shows the time comparison between my implementation and Psi4's.

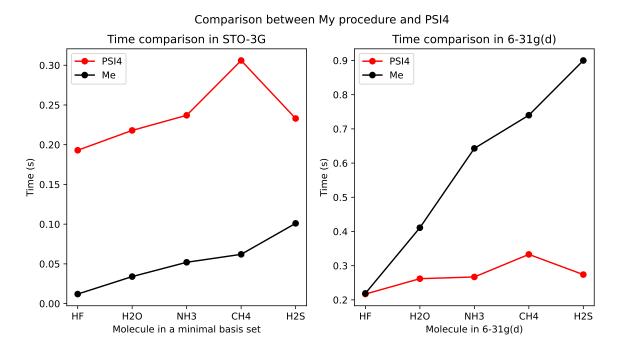


Figure 1:Comparison between my procedure and Psi4's

5 Conclusion

In conclusion, we have constructed the Self-Consistent-Field (SCF) procedure to carry out Hartree–Fock calculations. We started by proposing a set of basis functions and constructed its overlap matrix, then we used the overlap matrix to orthogonalize our entire Hartree-Fock–Roothan equation. Then we initialized a guess for our fock matrix to rebuild it again using the HFR equation. This procedure is variational, and we need to solve it iteratively until we meet a specific convergence condition.

My python implementation of the procedure was shown to be slightly faster in the minimal basis set STO-3G. I believe this is because Psi4 calculations need fractions of a second to start up before doing the calculations, which is about 0.2s. As we can see in the 6-31g(d) basis set, Psi4 started from 0.2s and was increased slightly for heavier and more complicated molecules. Whereas my calculations time increased dramatically.

References

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