

Firstly, this is my Mathematica code that I used for part 1&3:

```

 $\sigma_+ = \{\{0, 1\}, \{0, 0\}\}; \sigma_- = \{\{0, 0\}, \{1, 0\}\}$ 
 $c_{t,j} := \text{KroneckerProduct} @ N @ \text{ReplacePart}[\text{ReplacePart}[\text{Table}[\text{IdentityMatrix}[2], 6], \text{Table}[\{j-i\}, \{i, j-1\}] \rightarrow -\text{PauliMatrix}[3]], j \rightarrow \sigma_+]$ 
 $c_{j,-} := \text{KroneckerProduct} @ N @ \text{ReplacePart}[\text{ReplacePart}[\text{Table}[\text{IdentityMatrix}[2], 6], \text{Table}[\{j-i\}, \{i, j-1\}] \rightarrow -\text{PauliMatrix}[3]], j \rightarrow \sigma_-]$ 
 $\text{ListPlot}[\text{Sort} @ \text{Eigenvalues}\left[1 \left(\sum_{i=1}^5 (c_{t,i} c_{t,i+1}) + c_{t,6} c_{t,1}\right) + 2 \left(\sum_{i=1}^5 (c_{t,i} c_{t,i+1} c_{t,6} c_{t,i+1}) + c_{t,6} c_{t,1} c_{t,6} c_{t,1}\right)\right], \text{PlotStyle} \rightarrow \text{Black}, \text{PlotMarkers} \rightarrow \{\text{Automatic}, 3\}]$ 
 $\text{Grid}[\{\{c_5.c_5 + c_{t,5}.c_5 // \text{SparseArray}, c_5.c_{t,6} + c_{t,6}.c_5 // \text{SparseArray}\}, \{c_5.c_6 + c_6.c_5 // \text{SparseArray}, c_{t,5}.c_{t,6} + c_{t,6}.c_{t,5} // \text{SparseArray}\}\}]$ 

```

Figure 1: Mathematica code used for parts 1&3

1 Proving Fermionic Anti-Commutator Relations

For this part I did not do it analytically, I did it numerically. I set up a system with 7 particles, and defined the creation and destruction operators, then I applied the anti-commutator relations for two consecutive operators, $i, i + 1$. I found that the anti-commutator relations holds for this definition of operators, as seen in Figure. 2:

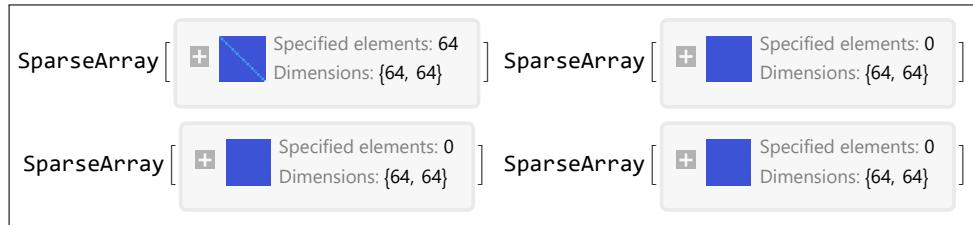


Figure 2: Results for $\{c_5, c_5^\dagger\}$; $\{c_5, c_6^\dagger\}$; $\{c_5, c_6\}$; $\{c_5^\dagger, c_6^\dagger\}$

2 Plotting Eigenvalues for a Hamiltonian

$$H = t \sum_i^6 c_i^\dagger c_{i+1} + U \sum_i^6 c_i^\dagger c_{i+1}^\dagger c_i c_{i+1}$$

Using the same system, these are the eigenvalues for above Hamiltonian:

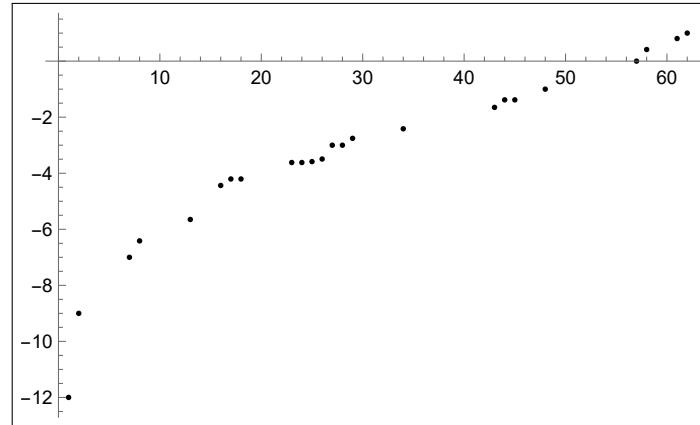


Figure 3: Sorted Eigenvalues of the Hamiltonian